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Zn-doping effects on the electrical resistivity of La_{1.85}Sr_{0.15}CuO_{4} under a magnetic field

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We measured the electrical resistivity of La_{1.85}Sr_{0.15}Cu_{1-y}Zn_{y}O_{4} (y=0,0.004,0.008,0.017) single crystals to study the magnetic-field dependence of the superconducting transitions. For y=0~0.008, as the magnetic field along the c axis is increased, the superconducting transition significantly broadens while the onset Tc remains unchanged in both the in-plane resistivity $\rho_{ab}$ and the out-of-plane resistivity $\rho_{c}$. This corresponds to the so-called broadening behavior. For y=0.017, the onset Tc is apparently reduced with increasing magnetic field, which we define as a parallel shift. This indicates that the in-plane superconducting coherence length drastically increases at y=0.017. We discuss these Zn-doping effects based on an inhomogeneous picture.

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I. INTRODUCTION

Early studies have shown that the electrical resistivity under a magnetic field exhibits a characteristic behavior in the high-$T_{c}$ superconductors: in the underdoped regime, the superconducting transition width $\Delta T_{c}$ broadens as the magnetic field along the c axis increases, while the onset $T_{c}$ remains almost unchanged.1–4 This is called a broadening behavior, which is now mainly ascribed to “superconducting fluctuations” and “vortex motion.” The superconducting fluctuations indicate that the superconducting order parameter fluctuates near $T_{c}$. The fluctuations become large in the high-$T_{c}$ superconductors because of their short superconducting coherence lengths $\xi$ and low dimensionalities.5 The vortex motion means that the Lorentz force drives vortices under a magnetic field so that the electric resistivity is induced.

Recently, it was reported that such a broadening behavior is eliminated for a certain kind of material; in La$_{2-x}$(Nd$_{x}$Sr$_{x}$Ba$_{x}$)$_{y}$CuO$_{4}$ around $x=1/8$, the onset $T_{c}$ is significantly decreased by a magnetic field, while $\Delta T_{c}$ remains constant.6–8 We refer to this behavior as a parallel shift. The materials, which exhibit the parallel shift, show a specific phenomenon; neutron scattering studies have revealed incommensurate (IC) orders due to spin correlations and most likely due to charge correlations as well. Such a parallel shift might be characteristic to a case where the IC static correlations appear. Although both the broadening behavior and the IC spin fluctuations are typical phenomena in the high-$T_{c}$ superconductors, the parallel shift has been observed only for the hole doping $x~1/8$. Therefore, a qualitatively different system is required to clarify whether the appearance of the IC static correlations and the parallel shift is intrinsically related or accidentally coincident.

In this paper, we report on the electrical resistivity measurements around $T_{c}$ under a magnetic field for La$_{1.85}$Sr$_{0.15}$Cu$_{1-y}$Zn$_{y}$O$_{4}$ (y=0,0.004,0.008,0.017) single crystals. It is known that impurity-free La$_{1.85}$Sr$_{0.15}$CuO$_{4}$ shows a “spin gap” in the antiferromagnetic (AF) IC spin excitations at low temperatures and that no static magnetic correlation have been observed for y=0.9–11 Recent studies have shown that the spin gap is gradually filled and that the IC static magnetic correlations appear by substituting Zn for Cu.12–14 The purpose of the present work is to investigate whether the resistivity curve near $T_{c}$ changes with the appearance of static magnetic correlations. We systematically measured both $\rho_{ab}(H\perp J_{ab})$ and $\rho_{c}(H\parallel c)$ for various Zn concentrations under a magnetic field along the c axis.

II. EXPERIMENTAL

Single crystals of La$_{1.85}$Sr$_{0.15}$Cu$_{1-y}$Zn$_{y}$O$_{4}$ (y=0,0.004,0.008,0.017) were grown using the traveling-solvent-floating-zone method. The concentrations of La, Sr, Cu, and Zn ions were precisely estimated by the inductively coupled plasma analysis. In addition, we measured the superconducting shielding signals with a superconducting quantum interference device magnetometer to determine bulk $T_{c}$. Each sample except for y=0 was cut from the same crystal rod as used in previous neutron scattering measurements,12–14 and the details of crystal growth and characterization are given elsewhere.14 As for y=0, the estimated Sr concentration is $x=0.142(2)$ and the shielding signal shows a sharp transition at $T_{c}$ (midpoint) of 37.4 K. The in-plane ($\rho_{ab}$) and out-of-plane ($\rho_{c}$) resistivities were measured by a standard four-probe dc method using a rotating sample holder to precisely align the samples against the magnetic field. We prepared long rectangular crystals, which sizes are typically $0.7\times0.7\times5$ mm$^{3}$ for $\rho_{ab}$ and $1\times1\times3$ mm$^{3}$ for $\rho_{c}$. After a silver paste was painted, the samples were heated under oxygen gas flow at 350–500 °C for 2 h to obtain a low contact resistance. The whole areas of both the ends were painted with a silver paste to assure a uniform current flow through the sample and gold wires were used as lead. After these processes, the contact resistance became less than 1 Ω. The samples were first cooled to the lowest temperature under magnetic field along the c axis up to 10 T. After the sample temperature was stabilized at each measuring point in a heating process, the resistivity measurements were repeated ten times and the data were averaged. Electric currents used for measurements were 10 mA for $\rho_{ab}$ and 1 mA for $\rho_{c}$ to ensure the accuracy for voltage reading. The Joule heating to the sample at low temperatures is esti-
La$_{1.85}$Sr$_{0.15}$Cu$_{1-y}$Zn$_y$O$_4$ with (a) $y=0$, (b) $y=0.004$, (c) $y=0.008$, and (d) $y=0.017$. Magnetic fields were applied along the $c$ axis from 0 to 10 T. The resistivity in the normal state $\rho_{\text{normal}}$ is estimated by the extrapolation from the linear $T$-dependent region just above $T_c$ at $H=0$ T. Dashed lines represent the estimated resistivity in the normal state. The inset shows the enlarged view of $\rho_{ab}$ for $y=0$. The arrow in the inset denotes the temperature at which a kink in the $\rho-T$ curve occurs.

**III. RESULTS**

Figures 1 and 2 show the temperature dependence of the in-plane resistivity $\rho_{ab}$ and the out-of-plane resistivity $\rho_c$ under magnetic fields along the $c$ axis. The $\rho_{ab}$ for $y=0$ and 0.008 in zero field show a two-step behavior, which probably comes from inhomogeneities of Sr or O in the sample. Since the transition at lower temperature approximately corresponds to the onset $T_c$ in the magnetization, it is considered that the transition at higher temperature occurs due to accidental connections of current paths of superconductivity and that the transition at lower temperature shows the bulk superconductivity. As for $y=0$, the resistivity exhibits a typical broadening behavior as previously reported,$^2$-$^4$ and the behavior in $\rho_{ab}$ is marked compared with that in $\rho_c$. The broadening of $\rho_{ab}$ is gradually suppressed by increasing Zn and $\rho_{ab}$ undergoes a parallel shift for $y=0.017$. In addition, we observed a kink in the resistivity curve of $\rho_{ab}$ for $y=0$ [the inset in Fig. 1(a)], and the kink gradually disappears with increasing Zn. However, no kink was observed in $\rho_c$ for all the samples.

In Fig. 3, we show magnetic-field dependences of two temperatures: one is the temperature at which the resistivity decreases to 95% of the resistivity in the normal state $\rho_{\text{normal}}(r=\rho/\rho_{\text{normal}}=0.95)$, $T_r(r=0.95)$, which is defined as the onset temperature, and the other is the temperature at $r=0.05$, $T_c(r=0.05)$, which corresponds to the end of transition. The difference between these two temperatures corre-
The transition width is apparently shifted to lower temperatures with increasing the magnetic field, indicating the parallel shift. The parallel shift behavior is marked in $\rho_{ab}$ compared with that in $\rho_c$. Furthermore, we note that $T_c(r=0.05)$ of $\rho_{ab}$ does not coincide with that of $\rho_c$ for $y=0$. Upon doping Zn, these two values become closer and finally coincide for $y=0.017$. The resistivity near $T_c(r=0.95)$ can be described by superconducting fluctuations, while the resistivity near $T_c(r=0.05)$ can be affected by vortex motion. We, however, cannot simply distinguish between the effect by superconducting fluctuations and that by vortex motion. We discuss the Zn-doping dependence of $T_c(r=0.05)$ and $T_c(r=0.95)$ in the Sec. IV.

IV. DISCUSSIONS

First, we discuss the magnetic-field dependence of $T_c(r=0.05)$ in $\rho_{ab}$ and $\rho_c$. As seen in Fig. 3, at higher fields, $T_c(r=0.05)$ in $\rho_{ab}$ does not coincide with that in $\rho_c$ for $y=0$. However, the difference becomes small upon doping Zn. Here we note that the gradient, $dH/dT$, becomes small as Zn increases in $\rho_{ab}$ but shows no Zn-doping dependence in $\rho_c$. It appears that there is an additional component in $\rho_{ab}$ and that the component is suppressed by Zn. Since the vortex motion originates from the Lorentz force, the resistivity contains the contribution of vortex motion only when the magnetic field is applied perpendicular to the current pass ($H \perp J$). In our measurement, $\rho_{ab}$ is contributed from the vortex motion because $J || ab$ is perpendicular to $H || c$. Furthermore, in the inset of Fig. 1(a), we see kinks in the $\rho-T$ curve at high magnetic fields. The previous study reported that the kink is attributed to a melting transition of the vortex lattice.\(^{15}\) Moreover, a recent study for the YBa$_2$Cu$_3$O$_{7-\delta}$ shows that a vortex lattice phase disappears with doping Zn,\(^{16}\) indicating the pinning of vortices by Zn. We speculate that the observed kink for $y = 0$ corresponds to a vortex liquid–vortex glass (or short-range vortex lattice) transition and that the vortex glass state is stabilized upon doping Zn.

Next, we observed the parallel shift behavior in $\rho_{ab}$ for $y=0.017$ unlike $y=0-0.008$, indicating that the superconducting fluctuations are small for $y=0.017$. The parallel shift was observed in particular materials, such as La$_{2-x}$(Nd,Sr,Ba)$_x$CuO$_4$ around $x=1/8$, which are expected to have static IC magnetic correlations or both charge and magnetic correlations.\(^{6-8}\) In our samples, the elastic IC magnetic signals were observed for $y=0.017$, while no signals appear for $y=0-0.008$.\(^{12-14}\) Thus, we conclude that the parallel shift is intrinsically related to the static IC correlations in the Zn-doped La$_{1.85}$Sr$_{0.15}$Cu$_{1-y}$Zn$_y$O$_4$ and that these phenomena are observed not only in the high-$T_c$ materials around the 1/8 hole concentration\(^{6-8}\) but also in impurity-doped systems. It indicates that the transport properties in the CuO$_2$ plane have some relationship to the magnetic one.

Let us discuss the Zn-doping dependence of superconducting coherence length $\xi$. According to the superconducting fluctuation theory with taking account of the $|\theta|^4$ term in the Ginzburg-Landau equation, the resistivity can be induced below the mean-field critical temperature $T_c(H)$.\(^{5}\) It was also reported that the onset temperature of diamagnetism was quite different from the $R=0$ temperature, which seems more
consistent with the onset temperature in resistivity in the high-$T_c$ cuprates.\textsuperscript{17} Thus, it is expected that $T_c(H)$ is located around the onset $T_c$. In the present analysis, we use $T_c(r=0.95)$ as $T_c(H)$. The critical field at $T=0$ K, $H_{c2}(0)$, is extrapolated using the Werthamer-Helfand-Hohenberg formula, given as\textsuperscript{18}

$$H_{c2}(0) = -0.7T_c \left( \frac{dH_{c2}(T)}{dT} \right)_{T=T_c}.$$

Then, $\xi$ is evaluated through the relation $H_{c2}(0) = \phi_0/2\pi\xi^2 = \phi_0/2\pi\xi_{\text{SC}}^2$ for $H\parallel c$, where $\phi_0$ is the magnetic flux quantum ($=\hbar c/2e = 206\,700\,\text{T}\,\text{Å}^2$). Figure 4(a) shows the Zn-doping dependence of $\xi_{ab}$, thus estimated. Note that $T_c(H)$ defined from $T_c(r=0.95)$ contains some uncertainties for $y=0-0.008$ showing a broad transition, which are affected by large superconducting fluctuations. This fact would lead to an overestimation of $\xi_{ab}$. To check the validity of our analysis, we also evaluated the coherence length $\xi_{ab}$ using $T_c(r=0.98)$ and $T_c(r=0.90)$ as $T_c(H)$. The estimated values of $\xi_{ab}$ are shown as error bars in Fig. 4(a). The lower error bars are virtually invisible because $\xi_{ab}$ estimated from $T_c(r=0.98)$ is almost same as that from $T_c(r=0.95)$. We have thus concluded that the uncertainties are small enough to discuss the Zn-doping dependence of $\xi_{ab}$ in further details.

As seen in Fig. 4(a), $\xi_{ab}$ for both $\rho_{ab}$ and $\rho_x$ demonstrate a tendency to increase upon doping Zn. It is remarkable that $\xi_{ab}$ for $\rho_{ab}$ is much longer than that for $\rho_x$ at $y=0.017$, reflecting the parallel shift behavior. Here $\xi_{ab}$ estimated from $\rho_{ab}$ is basically same as that from $\rho_x$. These results imply that $\xi_{ab}$ for $\rho_{ab}$ drastically changes at the boundaries of $y = 0.008$ and $y=0.017$, which we define as $y_0$. At the boundary $y_0$, the magnetic property also changes; there is no magnetic static correlation for $y < y_0$, but the IC magnetic peaks appear for $y > y_0$. It indicates that the transport properties in the CuO$_2$ plane may be related to the two-dimensional magnetic correlations. In YBa$_2$Cu$_3$O$_{6.9}$, Tomimoto et al.\textsuperscript{20} reported that the coherence length grows with increasing Zn and the increase in $\xi$, is larger than a theoretical prediction. The rapid increase in $\xi$ may give rise to a tendency toward the parallel shift, because the superconductivity comes close to three dimensional. The present results also show that $\xi_{ab}$ gradually increases except for $\rho_{ab}$ for $y=0.017$ and are consistent with their report for Zn-doped YBa$_2$Cu$_3$O$_{6.9}$. However, the sudden change in $\xi_{ab}$ for $\rho_{ab}$ at $y=y_0$ requires another scenario.

We offer an inhomogeneous model to explain the sudden growth of $\xi_{ab}$ for $\rho_{ab}$, at $y=y_0$. Our concept is schematically shown in Fig. 4(b). For the Zn-doped system, Nachumi et al.\textsuperscript{19} proposed an inhomogeneous picture, which they called the “swiss cheese” model; nonsuperconducting (non-SC) islands with the radii $\xi_{\text{SC}}$ around Zn ions reside in the superconducting (SC) sea. Furthermore, scanning tunneling microscopy (STM) studies for Zn-doped Bi$_2$Sr$_2$CaCu$_2$O$_{8+\delta}$ showed that Zn induces an intense quasiparticle-scattering resonance around Zn, indicating that the superconductivity is strongly suppressed within $\xi_{\text{SC}}$. The radius of an island, $\xi_{\text{SC}}$, estimated by $\mu$SR and STM correspond to 18 Å and 15 Å. Near $y=y_0$, the average distance between Zn ions $r_{\text{Zn-Zn}}$ roughly corresponds to $2\xi_{\text{SC}}$, and the IC magnetic peaks appear for $y > y_0$. From the inhomogeneous picture [see Fig. 4(b)], when $r_{\text{Zn-Zn}}$ reaches $2\xi_{\text{SC}}(y=y_0)$, the isolated nonsuperconducting islands become connected and the IC magnetic correlations can be stabilized. The non-SC islands float in the SC sea for $y < y_0$, while small SC lakes remain in the non-SC land for $y > y_0$. In the case of $y > y_0$, the SC regions are isolated and can be connected through a Josephson tunneling across the non-SC land, which can realize superconductivity. Some results suggest that the superconductivity near the critical concentration, $y_c$, at which $T_c$ is fully suppressed, is not a bulk property,\textsuperscript{22,23} supporting our idea.

Why is the superconducting fluctuation suppressed in the region of $y > y_0$? A possible reason can be that the three-dimensional superconductivity is realized in highly Zn-doped region. In fact, Tomimoto et al.\textsuperscript{20} reported that $\xi_{ab}$ increases upon Zn doping in YBa$_2$Cu$_3$O$_{6.9}$. However, such a crossover of dimensionality of superconductivity would not explain why $\xi_{ab}$ of $\rho_{ab}$ is different from that of $\rho_x$ for $y < y_0$. We speculate that $\rho_{ab}$ can be affected by a development of IC correlation in the two-dimensional CuO$_2$ plane. It is known that the static IC correlation competes with the
superconductivity. The development of IC correlation might interrupt the growth of superconductivity. To address this issue, further studies are required.

Another effect of Zn is that the resistivity exhibits an insulating behavior. It has been discussed that the insulating behavior is caused by a large residual scattering or a charge localization, which can destroy the superconductivity.

We performed electrical resistivity measurements under various magnetic fields for La$_{1.85}$Sr$_{0.15}$Cu$_{1-x}$Zn$_x$O$_4$ ($y=0,0.004,0.008,0.017$). A typical broadening behavior was observed for $y=0–0.008$, while the resistivity curves in $\rho_{ab}$ at $y=0.017$ exhibit a parallel shift; likewise, at $y=0.017$, the IC static correlation was observed unlike $y=0–0.008$. We conclude that the parallel shift is associated with the static IC correlation in the Zn-doped La$_{1.85}$Sr$_{0.15}$Cu$_4$O$_4$ and that these phenomena are observed not only in the high-$T_c$ superconductivity but also in impurity-doped system. The change from the broadening to the parallel shift suggests that the coherence length $\xi_{ab}$ for $y<\gamma_0$ (0.008 $<y<0.017$) is much shorter than that for $y>\gamma_0$. It can be explained by an inhomogeneous picture: when the average distance among Zn ions $r_{Zn-Zn}$ reaches the diameter of the non-SC region $2\xi_{SC}(y=\gamma_0)$, the non-SC areas connect with each other, while the remaining small SC regions are separated from one another. In the case, the IC correlation in a non-SC area develops and the superconductivity with small superconducting fluctuation is realized in the isolated SC regions.

**V. CONCLUSIONS**

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