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Zero-field (ZF) and longitudinal-field (LF) muon spin relaxation ($\mu$SR) measurements have been carried out in the overdoped single crystals of La$_{2-x}$Sr$_x$Cu$_{1-y}$Zn$_y$O$_4$ with $x=0.21$, $y=0$ and 0.01. In the case of the Zn-substituted crystal with $x=0.21$ and $y=0.01$, exhibiting no bulk superconductivity at low temperatures below 2 K. No clear coherent precession of the muon spin has been observed down to 20 mK. Slowing down of the Cu-spin fluctuations has also been observed in the non-Zn-substituted crystal with $x=0.21$ and $y=0$ at low temperatures below about 0.8 K in ZF, indicating that the observed magnetic anomaly is an intrinsic property at $x=0.21$. These $\mu$SR results suggest a strong possibility of the existence of a spin/charge ordered state at $x=0.21$ in the overdoped region.

The so-called stripe model of spins and holes has been suggested to explain a magnetically ordered state observed in La$_2$NiO$_4$, La$_2$NiO$_{4-\delta}$, and La$_2$Sr$_{1-x}$NiO$_4$. In these nickelates, static stripe order of Ni spins and doped holes appears at various values of the hole concentration, $p=\frac{2}{3}$, $\frac{1}{3}$, and $\frac{1}{2}$ per Ni. Following this result, the static stripe order of spins and holes was suggested in La$_{1.43}$Nd$_{0.57}$Sr$_{0.12}$CuO$_4$ around $p=\frac{1}{2}$ per Cu by the neutron-scattering experiment. Meanwhile, dynamical stripe correlations of spins and holes have been found to exist in a wide range of $p$ from the underdoped to the overdoped region in La$_{2-x}$Sr$_x$CuO$_4$ (LSCO) and tend to be statically stabilized around $p=\frac{1}{2}$ per Cu, which is well known as the so-called “stripe effect” (1-12). Therefore, on analogy with the case of the nickelates, a sort of stripe order of Cu spins and doped holes is expected to be stabilized in the overdoped region of LSCO.

Actually, Kakinuma et al. reported anomalous behavior of the transport phenomena with a little suppression of the superconductivity in Zn-substituted polycrystalline samples of La$_{2-x}$Sr$_x$Cu$_{1-y}$Zn$_y$O$_4$ (LSCO) around $x=0.22$, suggesting a possibility of the existence of an ordered state of spins and holes in the overdoped region. Moreover, Kawamata et al. recently developed the sample quality of LSCO and successfully grew a homogeneous single crystal with $x=0.21$ and $y=0.01$, exhibiting no bulk superconductivity at low temperatures at least above 2 K. They confirmed that bulk superconductivity appeared again as the Sr-doping level shifted to the lighter or heavier side from $x=0.21$ in the Zn-substituted single crystals with $y=0.01$. This fact means that the suppression of superconductivity in the Zn-substituted single crystal occurs singularly at $x=0.21$. Taking into account that the static stabilization of the stripe correlations of spins and holes in the high-$T_c$ cuprates seemed to compete with the superconductivity around $p=\frac{1}{2}$ per Cu in LSCO, they argued that the dynamical stripe correlations of spins and holes tended to be pinned by the substituted Zn in the Zn-substituted single crystal with $x=0.21$. That is to say, a stripe ordered state of spins and holes has been expected in the present single crystals of LSCO with $x=0.21$. Kawamata et al. also succeeded in growing a non-Zn-substituted single crystal with $x=0.21$ and $y=0$ which showed bulk superconductivity with a little suppressed superconducting transition temperature, $T_c$ ($T_c$ is about 25 K). This result suggested that the anomaly observed in LSCO around $x=0.21$ is an intrinsic property of the La-214 system.
FIG. 1. Zero-field μSR time spectra of the Zn-substituted single crystal of La$_{2-x}$Sr$_x$Cu$_{1-y}$Zn$_y$O$_4$ with $x = 0.21$ and $y = 0.01$ at various temperatures. Fast depolarization behavior of the muon spin is observed at low temperatures below 3.5 K, which is due to slowing down of the Cu-spin fluctuations. Solid lines below 3.5 K are the best fit results using the two-exponential function of $A_1e^{-\lambda_1t} + A_0e^{-\lambda_0t}$.

in the overdoped region. Therefore, we have carried out the muon spin relaxation (μSR) measurements in order to investigate the Cu-spin state in La$_{2-x}$Sr$_x$Cu$_{1-y}$Zn$_y$O$_4$ with $x = 0.21, y = 0$ and 0.01.

Single-crystal-growth procedures and transport properties of the grown single crystals of La$_{2-x}$Sr$_x$Cu$_{1-y}$Zn$_y$O$_4$ with $x = 0.21, y = 0$ and 0.01 are reported in a separate paper. Zero-field (ZF) and longitudinal-field (LF) μSR measurements were carried out at the RIKEN-RAL Muon Facility (RIKEN-RAL) at the Rutherford-Appleton Laboratory in the UK and at the Meson Science Laboratory at the Institute of Materials Structure Science in the High Energy Accelerator Research Organization (KEK-MSL) in Japan. A pulsed positive surface-muon beam with a momentum of 27 MeV/c was used in both facilities. Forward and backward counters were located on the upstream and downstream sides in the beam direction which was parallel to the initial muon-spin and LF directions. The asymmetry parameter, $A(t)$, was defined as $A(t) = [F(t) - \alpha B(t)]/[F(t) + \alpha B(t)]$, where $F(t)$ and $B(t)$ were total muon events counted by the forward and backward counters at a time $t$, respectively. The $\alpha$ is a calibration factor reflecting the relative counting efficiencies of the forward and backward counters. The initial asymmetry was defined as $A(0)$. The time evolution of the asymmetry parameter (the μSR time spectrum) was measured in ZF and LF. The single crystals were aligned as the CuO$_2$ planes were perpendicular to the initial muon-spin polarization. Low-temperature measurements using a $^3$He-$^4$He dilution refrigerator from 1 K down to 20 mK were carried out at KEK-MSL and high-temperature measurements above 0.3 K were carried out at RIKEN-RAL.

Figure 1 shows the ZF-μSR time spectra of the Zn-substituted single crystal obtained at various temperatures down to 0.3 K. The time spectrum shows a Gaussian type depolarization behavior above 5 K, indicating that only static nuclear-dipole fields exist at the muon site. The depolarization behavior deviates from the Gaussian-type at low temperatures below 3.5 K, indicating the appearance of a fast depolarizing component. At low temperatures below 2 K, the asymmetry of the muon spin polarization decreases quickly by nearly two thirds within a time range of 0.5 μsec after the arrival of muons in the sample, while the rest one third of the asymmetry depolarize slowly afterwards. The time spectrum has not changed at low temperatures between 0.3 K and 20 mK and clear coherent precession of the muon spin, which is similar to that observed in LSCO around $p = \frac{1}{8}$ per Cu, has not been observed even at 20 mK.

The time spectra at high temperatures above 5 K are analyzed using $A_0e^{-\lambda_0t}G_{z}(\Delta,t)$, where $\lambda_0$ is the depolarization rate and $A_0$ is the initial asymmetry. The term of $G_{z}(\Delta,t)$ is the Kubo-Toyabe function which represents the effect of nuclear-dipole fields distributed at the muon site with a distribution width of $\Delta$. After the fast depolarizing component appears below 3.5 K, the analysis using $G_{z}(\Delta,t)$ is no longer effective, so that we put $\Delta = 0$ making $G_{z}(\Delta,t) = 1$. Consequently, the two-exponential function of $A_1e^{-\lambda_1t} + A_0e^{-\lambda_0t}$ was used to describe simply the time spectra below 3.5 K, where $A_1$ and $\lambda_1$ are the initial asymmetry and depolarization rate of the fast depolarizing component, respectively. Accordingly, $A_0$ and $\lambda_0$ are regarded as the initial asymmetry and the depolarization rate of the slow depolarizing one, respectively.

Results of the analysis are summarized in Fig. 2. The $A_0$ decreases suddenly around 2 K with decreasing temperature for the Zn-substituted crystal, which correlates with the increase of $\lambda_1$. The $\lambda_0$ is enhanced with decreasing temperature and exhibits divergence at 2 K. These are typical characteristics of the case where a static magnetically ordered state of Cu spins appears in the La-214 systems as observed in the previous μSR studies.

The LF has been applied to investigate whether the Cu-spin state is dynamic or static at 0.3 K. Figure 3 shows the
variation with LF of the time spectrum of the Zn-substituted single crystal at low temperatures. The fast depolarization component is recovered with increasing LF and almost quenched by the application of LF of 1 K. In addition, it is found that the time spectrum in LF is nearly flat. These results mean that the observed fast depolarization of the muon spin is caused by a static internal field and that the Cu-spin fluctuations are almost suppressed at 0.3 K. Therefore, it is concluded that a static magnetically ordered state of Cu spins appears in the Zn-substituted single crystal at low temperatures. The magnetic transition temperature is determined from the results shown in Fig. 2 to be about 2 K. The absence of the coherent precession of the muon spin is corresponding to a disordered arrangement of the Cu spins like a spin-glass state.

It has been reported by Kawamata et al.\(^\text{14}\) that the resistivity of the present Zn-substituted single crystal with \(x = 0.21\) is metallic at high temperatures above about 80 K, but exhibits upturn below about 80 K in every direction of the crystal axes, showing an insulating behavior. In addition, no Meissner effect has been confirmed at temperatures above 2 K. These results suggest that doped holes are mobile at high temperatures above about 80 K but tend to be localized at lower temperatures. Since a static magnetically ordered state of Cu spins appears at low temperatures below 2 K in the same single crystal, it is suggested that the ground state of this Zn-substituted single crystal is a sort of coexisting state of statically ordered Cu spins and localized holes. Therefore, a plausible explanation of the observed disordered state can be that a stripe ordered state of spins and holes is pinned by the substituted Zn and stabilized being accompanied by the suppression of the Cu-spin fluctuations.\(^\text{11–14,16}\) As a result, it is concluded that the slowing down of the Cu-spin fluctuations is due to the Zn-substitution effect but to the intrinsic behavior at \(x = 0.21\) in LSCO. That is to say, the Cu-spin correlation is enhanced at \(x = 0.21\) and the Zn-substitution leads to a static order of Cu spins.

In summary, ZF- and LF-\(\mu\)SR measurements have been carried out in the single crystals of \(La_{2-x}SrCu_{1-y}Zn_0O_4\) with \(x = 0.21\) and \(y = 0\) and 0.01. Fast depolarization of the muon spin has been observed at low temperatures below 3.5 K in the Zn-substituted crystal and below 0.8 K in the non-Zn-substituted one. The difference is only the temperature at which the slowing down of the Cu-spin fluctuations starts to be observable by \(\mu\)SR. This difference can be reasonably explained by the model that the dynamical stripe correlations of spins and holes are pinned by the substituted Zn and stabilized being accompanied by the suppression of the Cu-spin fluctuations.\(^\text{11–14,16}\) As a result, it is concluded that the slowing-down behavior of the Cu-spin fluctuations is not due to the Zn-substitution effect but to the intrinsic behavior at \(x = 0.21\) in LSCO. That is to say, the Cu-spin correlation is enhanced at \(x = 0.21\) and the Zn-substitution leads to a static order of Cu spins.
substituted single-crystal. Taking into account the results of transport measurements, a coexisting state of statically ordered Cu spins and localized holes is expected to exist below 2 K in the Zn-substituted single crystal. A plausible explanation of this coexisting state is suggested to be that a stripe ordered state of spins and holes exists at low temperatures at $x=0.21$ in LSCO in the overdoped region also.

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